A FEM for the Square Root of the Laplace Operator

Enrique Otárola¹

¹Department of Mathematics University of Maryland, College Park.

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Outline

The Square Root of the Laplace Operator The Harmonic Extension and the Truncated Problem The Galerkin Approximation of the Harmonic Extension Numerical Implementation in deal.ii Numerical Results

Outline of Topics

The Square Root of the Laplace Operator

The Harmonic Extension and the Truncated Problem

The Galerkin Approximation of the Harmonic Extension

Numerical Implementation in deal.ii

Numerical Results



The Continuos Problem

The problem we shall be concerned with reads as follows: Given a smooth enough function f, find u such that

$$\begin{cases} (-\Delta)^{1/2} u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial\Omega, \end{cases}$$

where $\Omega \subset \mathbb{R}^d$, with d = 1, 2 is a bounded domain with a smooth boundary $\partial\Omega$ and $(-\Delta)^{1/2}$ denotes the square root of the Laplace operator $-\Delta$ in Ω with zero boundary values on $\partial\Omega$.





Concerning applications, nonlocal operators are of importance in a wide range of applications:

- Finance.
- Image Processing.
- Quasi-geostrophic flow models.
- Modeling hydraulic fractures and the evolution of a viscous liquid thin film.

The development of efficient computational solution techniques for this problem is fundamental.



Definition of the Square Root of the Laplacian

Spectral theory of the Laplacian $-\Delta$ in a smooth bounded domain Ω with zero Dirichlet boundary values. There exists a sequence of eigenvalues

$$0 < \lambda_1 < \lambda_2 \leq \cdots \leq \lambda_k \leq \cdots \to \infty$$

and,



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and, there exists an orthonormal basis $\{\varphi_k\}$ of $L^2(\Omega)$, where $\varphi_k \in H^1_0(\Omega)$ is an eigenfunction corresponding to λ_k :

$$\begin{cases} -\Delta \varphi_k = \lambda_k \varphi_k & \text{in } \Omega\\ \varphi_k = 0 & \text{on } \partial \Omega, \end{cases}$$
(1)

for $k = 1, 2, \cdots$. Regularity theory $\implies \varphi_k \in C^{\infty}(\bar{\Omega})$ for $k = 1, 2, \cdots$.



Definition of the Square Root of the Laplacian

The square root of the Dirichlet Laplacian, for a smooth function u, is given by

$$u = \sum_{k=1}^{\infty} c_k \varphi_k \mapsto (-\Delta)^{1/2} u = \sum_{k=1}^{\infty} c_k \lambda_k^{1/2} \varphi_k.$$

Density results $\implies (-\Delta)^{1/2} : H_0^1(\Omega) \to L^2(\Omega).$



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Density results $\implies (-\Delta)^{1/2} : H_0^1(\Omega) \to L^2(\Omega)$. Then if $f \in L^2(\Omega)$, we have $f = \sum_{i=1}^{\infty} f_{i}(\alpha_i) \implies \alpha_i = f_i \lambda_i^{-1/2}$

$$f = \sum_{k=1}^{n} f_k \varphi_k \implies c_k = f_k \lambda_k^{-1/2}$$

Numerical disadvantages: We need to find a sufficiently large number of eigenfunctions to obtain an accurate approximation.



Definition of the square root of the Laplacian

On the other hand, this operator can be seen as a singular integral

$$(-\Delta)^{1/2}u(x) = C_d \int_{\Omega} \frac{u(x) - u(y)}{|x - y|^{d+1}} dy,$$

where C_d is a normalization constant.



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Numerical disadvantages: the integrand is singular and the matrix obtained is dense. These inconveniences complicate the numerical computation.



Harmonic Extension

The approach presented by X. Cabre and J. Tan, (2010): relation between the nonlocal operator $(-\Delta)^{1/2}$ and its harmonic extension.

Given *u* defined in Ω , we consider its harmonic extension *v* in the cylinder $\mathcal{C} := \Omega \times (0, \infty)$, with *v* vanishing on $\partial_L \mathcal{C} := \partial \Omega \times [0, \infty)$.

$$\left\{ egin{array}{ll} -\Delta v &=& 0 \quad ext{in} \ \mathcal{C} = \Omega imes (0,\infty), \ v &=& 0 \quad ext{on} \ \partial_L \mathcal{C} = \partial \Omega imes [0,\infty), \ rac{\partial v}{\partial
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where ν is the unit outer normal to C at $\Omega \times \{0\}$. Then,

$$u = \operatorname{tr}_{\Omega} v := v(\cdot, 0)$$

Spaces for v and u

Space for v:

$$H^1_0(\mathcal{C}) := \{ v \in H^1(\mathcal{C}) | v = 0 \text{ a.e. on } \partial_L \mathcal{C} = \partial \Omega \times [0,\infty) \}.$$

Space for **u**:

$$\begin{aligned} \mathcal{V}_0(\Omega) &= H_{00}^{1/2}(\Omega) &= \left[H_0^1(\Omega), L^2(\Omega) \right]_{1/2,2} \\ &= \left\{ u \in H^{1/2}(\Omega) \middle| \int_{\Omega} \frac{u^2(x)}{d(x)} dx < +\infty \right\}. \end{aligned}$$



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Truncated Problem

Numerically, it cannot be solved because C is an infinite domain \implies We need to consider a suitable truncated problem.



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Why can we truncate the problem? Given M > 0, v satisfies

$$\|\nabla \mathbf{v}\|_{L^2(\Omega\times(M,\infty))}^2 < e^{-2\sqrt{\lambda_1}M} \|f\|_{\mathcal{V}_0(\Omega)^*}^2.$$

Consider *M* adequately large and define v^M in a bounded domain $C_M := \Omega \times (0, M)$, imposing a zero Dirichlet condition on $\Omega \times \{M\}$:

$$\begin{cases} -\Delta v^{M} = 0 & \text{in } \mathcal{C}_{M} = \Omega \times (0, M), \\ v^{M} = 0 & \text{on } \partial_{L} \mathcal{C}_{M} := \partial \Omega \times [0, M], \\ v^{M} = 0 & \text{on } \Omega \times \{M\}, \\ \frac{\partial v^{M}}{\partial \nu} = f & \text{on } \Omega \times \{0\}, \end{cases}$$



Weak Formulation of the Truncated Problem

Find $v^M \in H^1_0(\mathcal{C}_M)$ such that

$$\int_{\mathcal{C}_M} \nabla v^M \cdot \nabla \psi = \int_{\Omega} f \operatorname{tr}_{\Omega} \psi, \quad \text{for all } \psi \in H^1_0(\mathcal{C}_M).$$

$$\begin{aligned} H^1_0(\mathcal{C}_M) &:= \{ v \in H^1(\mathcal{C}_M) | v = 0 \text{ a.e. on } \partial_L \mathcal{C}_M, \\ \text{and } v = 0 \text{ a.e. on } \Omega \times \{M\} \}. \end{aligned}$$



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How good is this truncated problem?

$$M > \frac{1}{\sqrt{\lambda_1}} \ln \left(\frac{2}{\epsilon^2} \right) \implies \| v - v^M \|_{H^1_0(\mathcal{C}_M)} \le \epsilon \| f \|_{\mathcal{V}_0(\Omega)^*}.$$



Galerkin Approximation

Given a family of partitions T_k of the domain C_M into quadrilateral elements, we define for $n \ge 1$

$$\mathbb{V}^{n,0} := \{ v \in C^0(\overline{\mathcal{C}_M}) : v |_T \in \mathcal{Q}_n(T) \ \forall T \in \mathcal{T}_k \} \cap H^1_0(\mathcal{C}_M),$$

Galerkin approximation for v^M is given by: Find $v_h^M \in \mathbb{V}^{n,0}$ such that

$$\int_{\mathcal{C}_M} \nabla v_h^M \cdot \nabla w_h = \langle f, \operatorname{tr}_{\Omega} w_h \rangle, \quad \text{ for all } w_h \in \mathbb{V}^{n,0}.$$



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$$\int_{\mathcal{C}_M} \nabla \boldsymbol{v}_h^M \cdot \nabla \boldsymbol{w}_h = \langle f, \operatorname{tr}_{\Omega} \boldsymbol{w}_h \rangle, \quad \text{ for all } \boldsymbol{w}_h \in \mathbb{V}^{n,0}.$$

Standard FEM theory + truncated problem property implies

$$\|\boldsymbol{v}-\boldsymbol{v}_h^M\|_{H^1_0(\mathcal{C})} \leq C\left(\epsilon\|f\|_{\mathcal{V}_0(\Omega)^*} + h\|\boldsymbol{v}\|_{H^2(\mathcal{C}_M)}\right),$$

where $h = \max_{T \in T} h_T$.

Numerical Implementation in deal.ii

The main function is similar to the steps discussed in class. We implement Q_k adaptive refinement, Q_k global refinement, and Q_k exponential refinement, $k \ge 1$, in dimension d = 2, 3.



Numerical Implementation in deal.ii

Template class LaplaceProblem. The class that does all the work.

Member functions. They do what their names suggest.

```
private:
    void setup_system ();
    void assemble_system ();
    void solve ();
    void refine_grid ();
    void process_solution (const unsigned int cycle);
};
```



Numerical Implementation in deal.ii

Some member variables:

};

Triangulation<dim> DoFHandler<dim> FE_Q<dim> ConstraintMatrix SparsityPattern SparseMatrix<double> Vector<double> Vector<double> const RefinementMode ConvergenceTable triangulation; dof_handler; fe; hanging_node_constraints; sparsity_pattern; system_matrix; solution; system_rhs; refinement_mode; convergence table:



Numerical implementation in deal.ii

LaplaceProblem::run. The code is implemented in cycles. For each kind of mesh, we consider a fixed number of cycles and we solve.



Numerical Example

We consider $\Omega = (0, 1)$ and $f(x) = \pi sin(\pi x)$, then $C_M = (0, 1) \times (0, M)$ $u(x) = sin(\pi x)$ and $v(x, y) = sin(\pi x)e^{-\pi y}$.

$$\|v - v_h^M\|_{H^1_0(\mathcal{C})} \le C\left(\epsilon \|f\|_{\mathcal{V}_0(\Omega)^*} + h\|v\|_{H^2(\mathcal{C}_M)}\right), \quad \epsilon = \epsilon(M)$$



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M should change with *h* to get $\epsilon \approx h$

$$M = -\frac{2}{\pi} \ln \left(\frac{h}{\sqrt{2}} \right)$$



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Some Global Meshes



Figure: Degrees of freedom: 20, 81, 238 respectively.



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Results with Global Refinement



Results with Global Refinement

Computing L^2 and H^1 error norms.

```
Vector<float> difference_per_cell (triangulation.n_active_cells());
```

```
const double L2_error = difference_per_cell.l2_norm();
```



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Results with Global Refinement





Estimates for the Function u

What about u?



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Estimates for the Function u

What about u? Trace results imply an estimate for u

$$\begin{aligned} \|u - u_h^M\|_{H^{1/2}_{00}(\Omega)} &\leq \|v - v_h^M\|_{H^1_0(\mathcal{C})} \\ &\leq C\left(\epsilon \|f\|_{\mathcal{V}_0(\Omega)^*} + h\|v\|_{H^2(\mathcal{C}_M)}\right), \quad \epsilon = \epsilon(M) \end{aligned}$$

However, notice that this estimate is not optimal! Optimal estimate

$$\|u-u_h^M\|_{H^{1/2}_{00}(\Omega)} \leq C\left(\epsilon \|f\|_{\mathcal{V}_0(\Omega)^*} + h^{3/2}\|f\|_{H^1(\Omega)}\right), \quad \epsilon = \epsilon(M)$$



Results with Global Refinement



Figure: Decay of the L^2 , $H^{1/2}$ and H^1 norms of the error.



Exponential Refinement

We exploit the behavior of the real solution

$$v(x,y) = \sum c_k \varphi_k e^{-\sqrt{\lambda_k}y}, \text{ for all } (x,y) \in \mathcal{C},$$

to design an exponential mesh. 2D case: We do global refinement in x and exponential refinement in y. Using interpolation results we get

$$\|v - v_h^M\|_{H^1_0(\mathcal{C}_M)}^2 \leq C \sum_{k=1}^{N_y} (h_k^y)^3 e^{-\sqrt{\lambda_1}y_k} \leq CN^{-1},$$



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and finally we obtain

$$y_{k+1} = y_k + \frac{1}{k} N^{-2/3} e^{\sqrt{\lambda_1}/3y_k}.$$



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Some Exponential Meshes



Figure: Degrees of freedom: 54, 170, 627 respectively.



Numerical Results

Results with Exponential Refinement



Figure: Decay of the L^2 , $H^{1/2}$ and H^1 norms of the error.



Adaptive Refinement

The estimate

$$\|\boldsymbol{v}-\boldsymbol{v}_h^M\|_{H^1_0(\mathcal{C})} \leq C\left(\boldsymbol{\epsilon}\|f\|_{\mathcal{V}_0(\Omega)^*} + h\|\boldsymbol{v}\|_{H^2(\mathcal{C}_M)}\right),$$

is not computable and provides only asymptotic information. We create a mesh adapted to the function v. Basic ingredient:

$$\|\mathbf{v} - \mathbf{v}_h^M\|_{H^1_0(\mathcal{C})} \leq C_1 \mathcal{E}_{\mathcal{T}}(\mathbf{v}_h^M) \leq C_2 \left(\|\mathbf{v} - \mathbf{v}_h^M\|_{H^1_0(\mathcal{C})} + \operatorname{osc}_{\mathcal{T}}(\mathbf{v}_h^M)\right)$$



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Error Estimator Implemented in deal.ii

$$\mathcal{E}_{\mathcal{T}}^{2}(\mathbf{v}_{h}^{M},T)=rac{h_{T}}{24}\int_{\partial T}\left[rac{\partial \mathbf{v}_{h}^{M}}{\partial \mathbf{v}}
ight]$$



3D Numerical Example

We consider $\Omega = (0,1) \times (0,1)$ and $f(x) = \sqrt{2\pi} \sin(\pi x) \sin(\pi y)$, then $u(x) = \sin(\pi x) \sin(\pi y)$ and $v(x,y) = \sin(\pi x) \sin(\pi y) e^{-\sqrt{2\pi}y}$.

We have optimal estimates for every refinement: adaptive, exponential and global. We show the results obtained using Adaptivity.



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An Adaptive Mesh. M = 4.







Figure: Degrees of freedom: 13435.



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Convergence Table for v

n cells		H ¹ -error		L ² -error	
0	4	4.016e-01	-	4.790e-02	-
1	32	6.419e-01	-0.68	7.156e-02	-0.58
2	228	6.252e-01	0.04	6.094e-02	0.23
3	1628	4.190e-01	0.58	2.845e-02	1.10
4	11400	2.312e-01	0.86	8.959e-03	1.67
5	79265	1.188e-01	0.96	2.394e-03	1.90
6	549238	5.983e-02	0.99	6.091e-04	1.97



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Results with Adaptive Refinement



Figure: Decay of the L^2 , $H^{1/2}$ and H^1 norms of the error.



Results with Adaptive Refinement



Figure: u_h^M and v_h^M with 13435 degrees of freedom.



3D Numerical Example

We consider the following numerical example. Given a smooth function $f(x, y) = \sqrt{2\pi} \sin(\pi x) \sin(\pi y)$, find *u* such that

$$\begin{cases} (-\Delta)^{1/2} u = f & \text{in } \Omega, \\ u = 0 & \text{on } \partial \Omega, \end{cases}$$

where $\Omega=(-1,1)^2$ - $\operatorname{disk}_{(0,0)}(0.5).$



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Figure: Meshes for z = 0 and z = 4.



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An Adaptive Mesh. M = 4.







An Adaptive Mesh. M = 4.



Figure: u_h^M computed with 22492 degrees of freedom.



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Questions?



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Image: A (1)